BSE in



http://www.exciting-code.org

Stephan Sagmeister, University of Leoben, Austria

Outline

- Review BSE
- Implementation within LAPW and exciting
 - Elements of the BSE
 - Direct interaction term
 - Exchange interaction term
 - Screened Coulomb interaction
 - q-dependent screening
 - Singularities
- TDDFT kernel from BSE
- Parallelization

BSE

 e-h eigenvalue problem (effective electron-hole Hamiltonian)

$$\sum_{v'c'\mathbf{k'}} H_{vc\mathbf{k},v'c'\mathbf{k'}} A_{v'c'\mathbf{k'}}^{\lambda} = \mathcal{E}^{\lambda} A_{vc\mathbf{k}}^{\lambda}$$

macroscopic dielectric function including excitonic effects

$$\epsilon_{\mathbf{M}}^{\alpha\alpha}(\omega) = 1 - \frac{8\pi}{V} \sum_{\lambda} |t_{\lambda}^{\alpha}|^2 \left[\frac{1}{\omega - \mathcal{E}_{\lambda} + i\delta} + \frac{1}{-\omega - \mathcal{E}_{\lambda} - i\delta} \right]$$

$$t_{\lambda}^{\alpha} = \sum_{vc\mathbf{k}} A_{vc\mathbf{k}}^{\lambda} \underbrace{\left\langle v\mathbf{k} | \hat{p}_{\alpha} | c\mathbf{k} \right\rangle}_{\varepsilon_{c\mathbf{k}} - \varepsilon_{v\mathbf{k}}}$$

The BSE Hamiltonian

Three parts

$$H^{\text{eff}} = H^{\text{diag}} + H^{x} + H^{\text{dir}}$$

Repulsive exchange interaction

$$H_{vc\mathbf{k},v'c'\mathbf{k}'}^{x} = \int d^{3}r d^{3}r' \psi_{v\mathbf{k}}(\mathbf{r}) \psi_{c\mathbf{k}}^{*}(\mathbf{r}) \bar{v}(\mathbf{r},\mathbf{r}') \psi_{v'\mathbf{k}'}^{*}(\mathbf{r}') \psi_{c'\mathbf{k}'}(\mathbf{r}')$$

Attractive direct interaction

$$-W$$

$$H_{vc\mathbf{k},v'c'\mathbf{k'}}^{dir} = -\int d^3r d^3r' \psi_{v\mathbf{k}}(\mathbf{r}) \psi_{c\mathbf{k}}^*(\mathbf{r'}) W(\mathbf{r},\mathbf{r'}) \psi_{v'\mathbf{k'}}^*(\mathbf{r}) \psi_{c'\mathbf{k'}}(\mathbf{r'})$$

Diagonal part

$$H_{vc\mathbf{k},v'c'\mathbf{k'}}^{\text{diag}} = (E_{c\mathbf{k}} - E_{v\mathbf{k}})\delta_{vv'}\delta_{cc'}\delta_{\mathbf{k}\mathbf{k'}}$$



The BSE Hamiltonian

- H(+) resonant part
- H(–) anti-resonant part
- H(#), H(#') coupling parts (beyond Tamm-Dancoff approximation, not yet implemented)

$$H^{\text{eff}} = \begin{pmatrix} H_{v_1c_2,v_3c_4}^{(+)} & H_{v_1c_2,c_3v_4}^{(\#)} \\ H_{c_1v_2,v_3c_4}^{(\#')} & H_{c_1v_2,c_3v_4}^{(-)} \end{pmatrix}$$



The BSE Hamiltonian

Effective e-h Hamiltonian

$$H^{\text{eff}} = H^{\text{diag}} + H^{x} + H^{\text{dir}}$$

• exchange part $\langle vc|_{6}^{5} \sim \sqrt[7]{v'c'} > H_{vc\mathbf{k},v'c'\mathbf{k}'}^{x} = \frac{1}{V} \sum_{\mathbf{G}} M_{vc\mathbf{k}}^{*}(\mathbf{G}) \bar{v}_{\mathbf{G}}(0) M_{v'c'\mathbf{k}'}(\mathbf{G})$



Common ingredients

 Exchange and direct part of BSE Hamiltonian can be written in terms of:

$$M_{nn'\mathbf{k}}(\mathbf{q} + \mathbf{G}) = \langle n\mathbf{k}|e^{-i(\mathbf{q}+\mathbf{G})\mathbf{r}}|n'\mathbf{k} + \mathbf{q}\rangle$$

 $P_{nn'\mathbf{k}}^{j} = \langle n\mathbf{k}|-i\nabla_{j}|n'\mathbf{k}\rangle$

wave function within APW-like basis sets

Atomic like basis functions in muffin-tin: spherical harmonics expansion

$$\phi_{n\mathbf{k}}^{\mathrm{MT}}(\mathbf{r}) = \sum_{\alpha lmp} A_{lmp}^{\alpha}(n\mathbf{k}) u_{lp}^{\alpha}(r_{\alpha}) Y_{lm}(\hat{\mathbf{r}}_{\alpha})$$

$$\phi_{n\mathbf{k}}^{\mathbf{I}}(\mathbf{r}) = \sum_{\mathbf{G}} c_{n\mathbf{k}+\mathbf{G}} e^{i(\mathbf{k}+\mathbf{G})\mathbf{r}}$$

$$A_{lmp}^{\alpha}(n\mathbf{k}) = \sum_{\mathbf{G}} c_{n\mathbf{k}+\mathbf{G}} A_{lmp}^{\alpha}(\mathbf{k}+\mathbf{G})$$

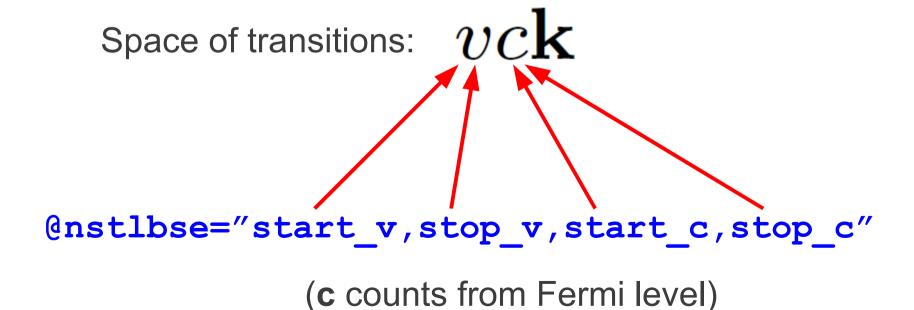
Plane waves for interstitial region, slow changes in wave function

 α ... atom

n ... band index

Controlling the Hamiltonian

$$H_{vc\mathbf{k},v'c'\mathbf{k'}}$$



Screening

Make use of TDDFT response at ω=0

$$\epsilon_{\mathbf{G}\mathbf{G}'}(\mathbf{q},\omega) = \delta_{\mathbf{G}\mathbf{G}'} - \chi_{\mathbf{G}\mathbf{G}'}^{0}(\mathbf{q},\omega)v_{\mathbf{G}'}(\mathbf{q})$$

$$\chi_{\mathbf{GG'}}^{\mathrm{KS}}(\mathbf{q},\omega) = \frac{1}{V} \sum_{nm\mathbf{k}} \frac{f_{n\mathbf{k}} - f_{m\mathbf{k}+\mathbf{q}}}{\varepsilon_{n\mathbf{k}} - \varepsilon_{m\mathbf{k}+\mathbf{q}} + \omega + i\delta} M_{nm\mathbf{k}}(\mathbf{q},\mathbf{G}) M_{nm\mathbf{k}}^*(\mathbf{q},\mathbf{G'})$$

Statically screened Coulomb interaction

$$W_{\mathbf{GG'}}(\mathbf{q}) = \epsilon_{\mathbf{GG'}}^{-1}(\mathbf{q})v_{\mathbf{G'}}(\mathbf{q})$$



Screened interaction at Gamma-point

• Subcell average $W_{\mathbf{GG'}}'(\mathbf{q}) = \frac{1}{V_q} \int_{V_\mathbf{q}} \mathrm{d}^3 p W_{\mathbf{GG'}}(\mathbf{p})$ on spherical Lebedev-Laikov grids (Freysoldt, CPC 2007)

$$\int_{V_{\mathbf{q}}} d^3 p f(\mathbf{p}) = \int d\hat{\Omega}_{\mathbf{p}} \int dp \lambda_{\mathbf{q}}(p, \hat{\Omega}_{\mathbf{p}}) f(\mathbf{q} + \mathbf{p})$$

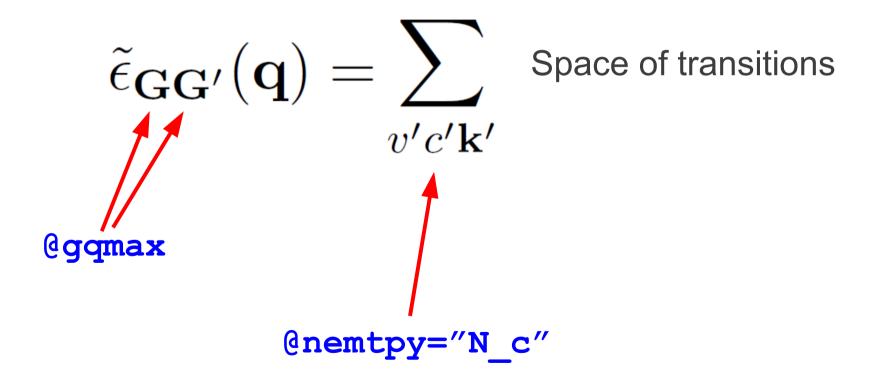
$$\bullet \ \ \text{Head of W:} \quad W_{00}'(0) = \frac{1}{V_q} \int \mathrm{d}\hat{\Omega}_{\mathbf{p}} \epsilon_{00}^{-1}(\hat{\Omega}_{\mathbf{p}}) \int \mathrm{d}p \lambda_{\Gamma}(p,\hat{\Omega}_{\mathbf{p}}) \frac{4\pi^2}{p^2}$$

• Approximation for head $\frac{1}{V_q} \int \mathrm{d}^3 p \, \frac{4\pi^2}{q^2} = \frac{2}{\pi} V q_s$

Effective subcell radius: q_s



Controlling the screening





BSE for core-levels

Trick: Treat core states as local orbitals valence states. More information:

- Claudia's talk (tomorrow)
- Weine's tutorial (tomorrow)



TDDFT with a xc-kernel derived from BSE

MBPT xc kernel

 TDDFT kernel derived from BSE (Nanoquanta groups), excitonic part, approximation to 1st order in W:

$$\int_{\mathbf{6}}^{\mathbf{7}} H_{vc\mathbf{k},v'c'\mathbf{k'}}^{dir} = -\int d^3r d^3r' \psi_{v\mathbf{k}}(\mathbf{r}) \psi_{c\mathbf{k}}^*(\mathbf{r}') W(\mathbf{r},\mathbf{r}') \psi_{v'\mathbf{k'}}^*(\mathbf{r}) \psi_{c'\mathbf{k'}}(\mathbf{r}')$$

Yields to 1st order in W:

$$f_{\rm xc}^{\rm e-h} pprox \chi_{\rm QP}^{-1} GGWGG\chi_{\rm QP}^{-1}$$

MBPT xc kernel

Reformulation of TDDFT Dyson equation (Sottile, PRL 2003)

$$\begin{split} \chi_{\mathbf{G}\mathbf{G'}}(\omega) &= \sum_{\mathbf{G}_1\mathbf{G}_2} \chi_{\mathbf{G}\mathbf{G}_1}^0(\omega) \left[\chi_{\mathbf{G}_1\mathbf{G}_2}^0(\omega) - \sum_{\mathbf{G}_3} \chi_{\mathbf{G}_1\mathbf{G}_3}^0 v_{\mathbf{G}_3} \chi_{\mathbf{G}_3\mathbf{G}_2}^0 - T_{\mathbf{G}_1\mathbf{G}_2}^{\mathrm{xc}} \right]^{-1} \\ &\times \chi_{\mathbf{G}_2\mathbf{G'}}^0. \end{split}$$

Reformulation of xc kernel part

$$T_{\mathbf{GG'}}^{\mathrm{xc}}(\omega) = \sum_{vc\mathbf{k}} \sum_{v'c'\mathbf{k}'} \frac{M_{vc\mathbf{k}}(\mathbf{G})W_{vc\mathbf{k},v'c'\mathbf{k}'}M_{v'c'\mathbf{k}'}^{*}(\mathbf{G}')}{(\varepsilon_{v\mathbf{k}} - \varepsilon_{c\mathbf{k}} + \omega + i\delta)(\varepsilon_{v'\mathbf{k}'} - \varepsilon_{c'\mathbf{k}'} + \omega + i\delta)}$$

$$T_{\rm xc} = \chi^0 f_{\rm xc} \chi^0$$

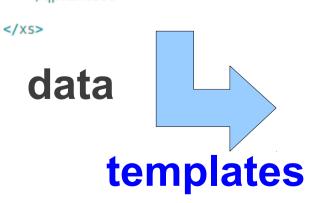
Important parameters

(input/xs/)

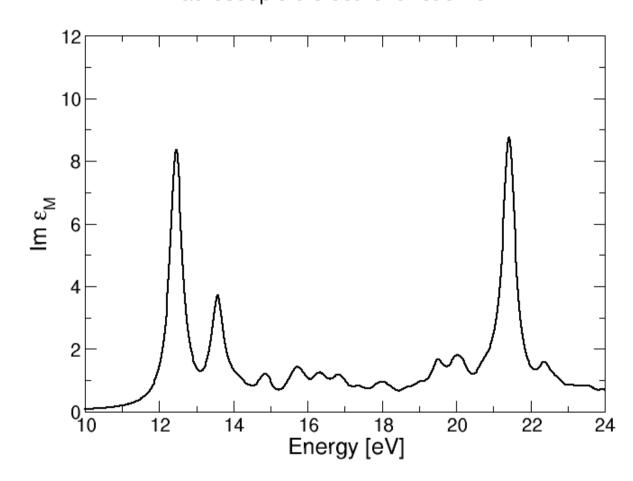
- @ngridk, @ngridq k/q-point grid sizes
- @rgkmax |G+k|_{max}R_{min} basis set size for wavefunction
- screening/@nempty #unocc. states for screening
- @gqmax |G+q|_{max} local field effects size
- energywindow range of energies for spectrum
- bse/@nstlbse ranges for occupied and unocc. states in BSE Hamiltonian
- bse/@bsetype type of BSE Hamiltonian (RPA, singlet, triplet)



In practice:



Macroscopic dielectric function of LiF



Parallelization

MPI parallelization

- Wavefunctions: #k-points
- Matrix elements #k-points
- screening #q-points
- Exchange interaction #k-points
- Direct interaction #(kk')-pairs
- Diagonalization not yet

Data parallelism throughout – no side effects



Now grab your coffee and we meet at the excercises...

